

The Local Temperature Dependence of Fluorescent Centres in PVA Films on the Excitation Wavenumber

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The temperature difference $\Delta T = T^* - T$ (where T^* is the local temperature and T the ambient temperature) is studied as a function of the excitation wavenumber $\tilde{\nu}_{\text{exc}}$ in the longwave region of the absorption band of 4-amino-4'-nitrostilbene (ANS) and 4-dimethylamino-4'-nitrostilbene (DNS) as well as rhodamine S, rhodamine 6G and Na-fluorescein in poly(vinyl alcohol) (PVA) films heated up to 403 K. Local temperatures T^* are determined by a universal Kennard-Stepanov relation. A linear dependence of ΔT on $\tilde{\nu}_{\text{exc}}$ has been found for all investigated luminescent compounds. For dyes, due to the big spectral overlaps between the absorption and fluorescence bands, it was possible to excite also in the anti-Stokes region. In this case for $\Delta T = 0$, i.e. when $T^* = T$, one has $\tilde{\nu}_{\text{exc}} = \tilde{\nu}'_{0-0}$. In view of the slow geometrical relaxation of excitation energy excess in PVA polymers, the relaxation time τ_R for thermal equilibration between the luminescent molecule and the vicinal polymer exceeds distinctly the mean fluorescence lifetime τ_F .

Key words: Kennard-Stepanov Universal Relation; Local Temperature; PVA Fluorescent Films.

1. Introduction

In [1], local temperatures T^* have been determined based on a universal Kennard-Stepanov relation [2, 3] for a series of luminescent molecules in unheated and heated PVA films. For unheated and heated films of the same luminescent compound different values of T^* were obtained. Reproducible results can only be obtained for rigidified films heated above the glassy transition temperature $T = 358$ K. In a freshly prepared unheated PVA film there exist regions of a different degree of polymerization, and such an unrigidified film forms an inhomogeneous environment.

For the luminescent molecules studied in [1], in the PVA polymer heated up to 403 K a local heating takes place which is characterized by a temperature T^* distinctly higher than the ambient temperature T . Higher temperature T^* evidences excitation energy excess, and this is the reason of the observed subtle dependence of fluorescence band peak location on the excitation wavelength in the longwave absorption band. A particularly strong shift of the fluorescence band is observed for the strongly polar molecules 4-amino-4'-nitrostilbene (ANS) and 4-dimethylamino-4'-nitrostilbene (DNS) in heated PVA films.

Therefore it is necessary to study the difference $\Delta T = T^* - T$ as a function of the excitation wavenumber $\tilde{\nu}_{\text{exc}}$ in the longwave absorption band. In the present paper

the following molecules are studied: ANS, DNS, rhodamine S, rhodamine 6G, and Na-fluorescein only in heated up to 403 K PVA films.

2. Experimental

Isotropic poly(vinyl alcohol) films (PVA) were prepared by the method described in [4, 5]. Absorption and fluorescence spectra were measured with the apparatus described in [6, 7].

3. Results and Discussion

The temperatures T^* were determined by the universal Kennard-Stepanov relation between the extinction coefficient $\varepsilon(\tilde{\nu})$ and the fluorescence intensity $f(\tilde{\nu})$ [2, 3], written in the modified form [8]

$$F(\tilde{\nu}, T^*) \equiv \ln \frac{f(\tilde{\nu}, T^*)}{\tilde{\nu}^3 \cdot \varepsilon(\tilde{\nu}, T^*)} = -\frac{hc\tilde{\nu}}{kT^*} + \text{const}, \quad (1)$$

where h and k are the Planck and Boltzmann constants, respectively, and c is the velocity of light.

Figures 1 and 2 present experimental data for ANS and DNS in heated PVA films for three selected wavenumbers of the exciting light $\tilde{\nu}_{\text{exc}}$. A distinct dependence of the fluorescence band location on $\tilde{\nu}_{\text{exc}}$ is observed. The

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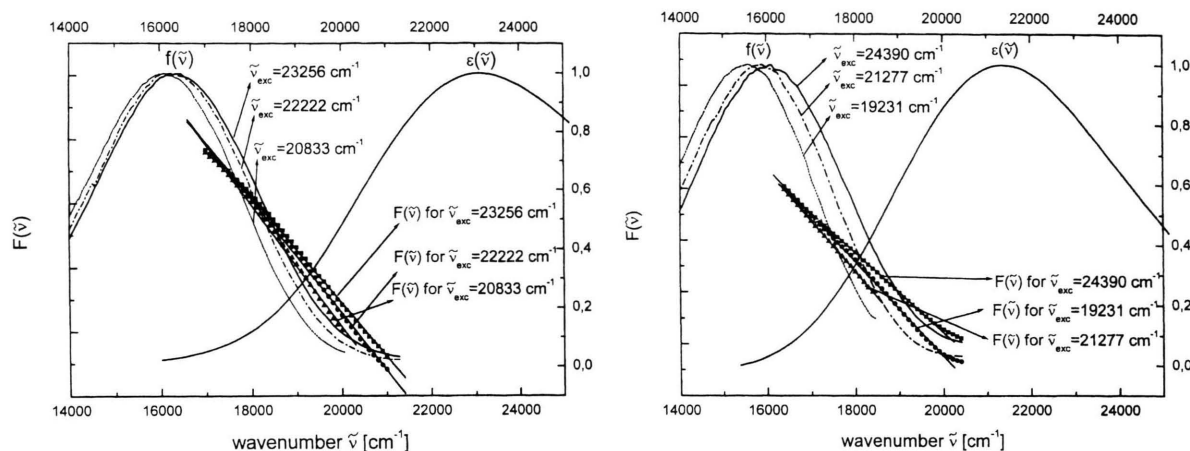


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Figs. 1 and 2. Absorption $\varepsilon(\tilde{\nu})$, fluorescence, $f(\tilde{\nu})$, spectra and the corresponding $F(\tilde{\nu})$ functions at different excitation wave-numbers, $\tilde{\nu}_{\text{exc}}$, for ANS (Fig. 1) and DNS (Fig. 2) in PVA films.

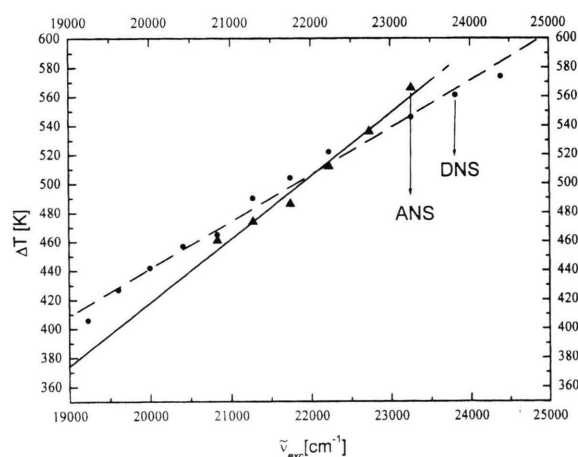


Fig. 3. The increments of temperatures, $\Delta T = T^* - T$, at different excitation wavenumbers $\tilde{\nu}_{\text{exc}}$ for ANS and DNS in PVA films.

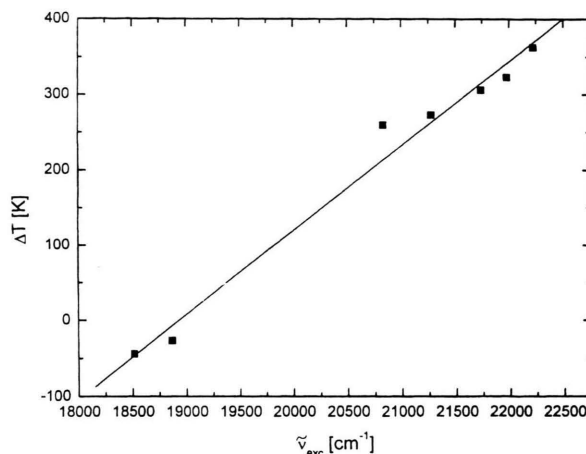
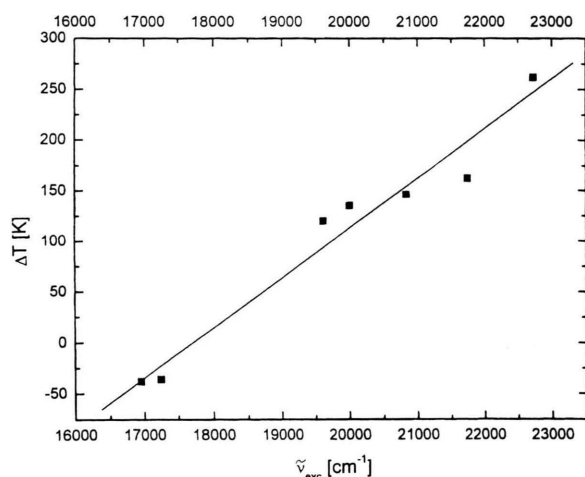
temperatures T^* listed in Tables 1 and 2, based on the linear dependence of $F(\tilde{\nu})$ given by (1), exceed for different $\tilde{\nu}_{\text{exc}}$ the ambient temperature T similarly as for complex molecules in gaseous phase [9].

For $\tilde{\nu}_{\text{exc}} > \tilde{\nu}'_{0-0}$ the temperature difference is positive, $\Delta T = T^* - T > 0$, which means that excited molecules possess an excess of vibrational energy compared to the equilibrated energy characterized by the ambient temperature T . Figure 3 shows the dependence of ΔT on $\tilde{\nu}_{\text{exc}}$ for ANS and DNS. Similar linear dependencies have been found for rhodamine S, rhodamine 6G, and Na-fluorescein in heated PVA films (Figs. 4 and 5). In the latter cas-

Table 1. Excited state local temperatures of stilbene derivatives in PVA films determined for different wavenumbers of exciting light.

| Compound | $\tilde{\nu}_{\text{exc}}$ [cm ⁻¹] | $\tilde{\nu}'_{0-0}$ [cm ⁻¹] | $\Delta \tilde{\nu}$ [cm ⁻¹] | T^* [K] | ΔT [K] |
|----------|---|---|---|--------------|-------------------|
| ANS | 23256 | 19130 | 4126 | 862 | 566 |
| | 22727 | 19080 | 3647 | 832 | 536 |
| | 22222 | 19020 | 3202 | 808 | 512 |
| | 21739 | 18920 | 2819 | 782 | 486 |
| | 21277 | 18890 | 2387 | 770 | 474 |
| | 20833 | 18820 | 2013 | 757 | 461 |
| DNS | 24390 | 18384 | 6006 | 856 | 560 |
| | 23809 | 18335 | 5474 | 836 | 540 |
| | 23255 | 18271 | 4984 | 813 | 517 |
| | 22727 | 18238 | 4489 | 795 | 499 |
| | 22222 | 18206 | 4016 | 777 | 481 |
| | 21739 | 18173 | 3566 | 761 | 465 |
| | 21276 | 18109 | 3167 | 753 | 457 |
| | 20833 | 18070 | 2763 | 715 | 419 |
| | 20408 | 18034 | 2374 | 705 | 409 |
| | 20000 | 18000 | 2000 | 682 | 386 |
| | 19607 | 17930 | 1677 | 657 | 361 |
| | 19230 | 17860 | 1370 | 625 | 329 |

es the samples could also be excited in the anti-Stokes region due to the large overlap between the absorption and fluorescence spectra. For the anti-Stokes region, the reverse relation $\tilde{\nu}_{\text{exc}} < \tilde{\nu}'_{0-0}$ holds and $\Delta T \leq 0$, which means that the supply of vibrational energy of excited molecules is smaller or equal to the equilibrated energy corresponding to T . From the Table 2 and Figs. 4 and 5 it is seen that for $\Delta T = 0$, which means that temperature T^* of the excited molecules is equal to the ambient temperature T , the relation $\tilde{\nu}_{\text{exc}} = \tilde{\nu}'_{0-0}$ holds true. Tables 1



Figs. 4 and 5. The increments of temperatures, $\Delta T = T^* - T$, at different excitation wavenumbers $\tilde{\nu}_{\text{exc}}$ for rhodamine S (Fig. 4) and Na-fluorescein (Fig. 5) in PVA films.

Table 2. Excited state local temperature of dye molecules in PVA films determined for different wavenumbers of exciting light.

| Compound | $\tilde{\nu}_{\text{exc}}$ [cm ⁻¹] | $\tilde{\nu}_{0-0}$ [cm ⁻¹] | $\Delta\tilde{\nu}$ [cm ⁻¹] | T^* [K] | ΔT [K] |
|----------------|---|--|--|--------------|-------------------|
| Rhodamine S | 22727 | 17844 | 4883 | 558 | 262 |
| | 21739 | 17861 | 3878 | 458 | 162 |
| | 20833 | 17871 | 2962 | 442 | 146 |
| | 20000 | 17867 | 2133 | 431 | 135 |
| | 19608 | 17844 | 1764 | 416 | 120 |
| | 17241 | 17738 | -497 | 260 | -36 |
| | 16949 | 17711 | -762 | 258 | -38 |
| Rhodamine 6G | 22222 | 18242 | 3980 | 493 | 197 |
| | 20833 | 18242 | 2591 | 469 | 173 |
| | 20000 | 18242 | 1758 | 450 | 154 |
| | 19607 | 18239 | 1368 | 430 | 134 |
| | 17241 | 18138 | -897 | 288 | -8 |
| Na-fluorescein | 16949 | 18147 | -1198 | 267 | -29 |
| | 22222 | 19567 | 2655 | 659 | 363 |
| | 21978 | 19576 | 2402 | 619 | 323 |
| | 21739 | 19550 | 2189 | 602 | 306 |
| | 21277 | 19567 | 1710 | 569 | 273 |
| | 20833 | 19556 | 1277 | 556 | 260 |
| | 18868 | 19392 | -524 | 269 | -27 |
| | 18519 | 19269 | -750 | 252 | -44 |

Table 3. Values of wavenumbers $\tilde{\nu}'_{0-0}$ obtained from interpolation of experimental dependencies $\Delta T(\tilde{\nu}_{\text{exc}})$ for $\Delta T = 0$ K.

| Compound | | |
|------------------------|------------------------|------------------------|
| Rhodamine 6G | Rhodamine S | Na-fluorescein |
| 17170 cm ⁻¹ | 17700 cm ⁻¹ | 18935 cm ⁻¹ |

and 2 show apart from T^* the values of $\tilde{\nu}_{0-0}$ obtained from the location of absorption and fluorescence peaks for different $\tilde{\nu}_{\text{exc}}$. The best agreement between the values of $\tilde{\nu}_{0-0}$ (Table 2) and $\tilde{\nu}'_{0-0}$ (Table 3) was obtained for rhodamine S.

The excitation energy excess for $\tilde{\nu}_{\text{exc}} > \tilde{\nu}_{0-0}$ causes in the rigid PVA polymer local heating during the emission process. Having in mind the slow geometrical relaxation of this energy excess, the relaxation time, τ_R , for thermal equilibration between the luminescent molecule and the surrounding polymer exceeds distinctly the mean fluorescence lifetime τ_F .

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